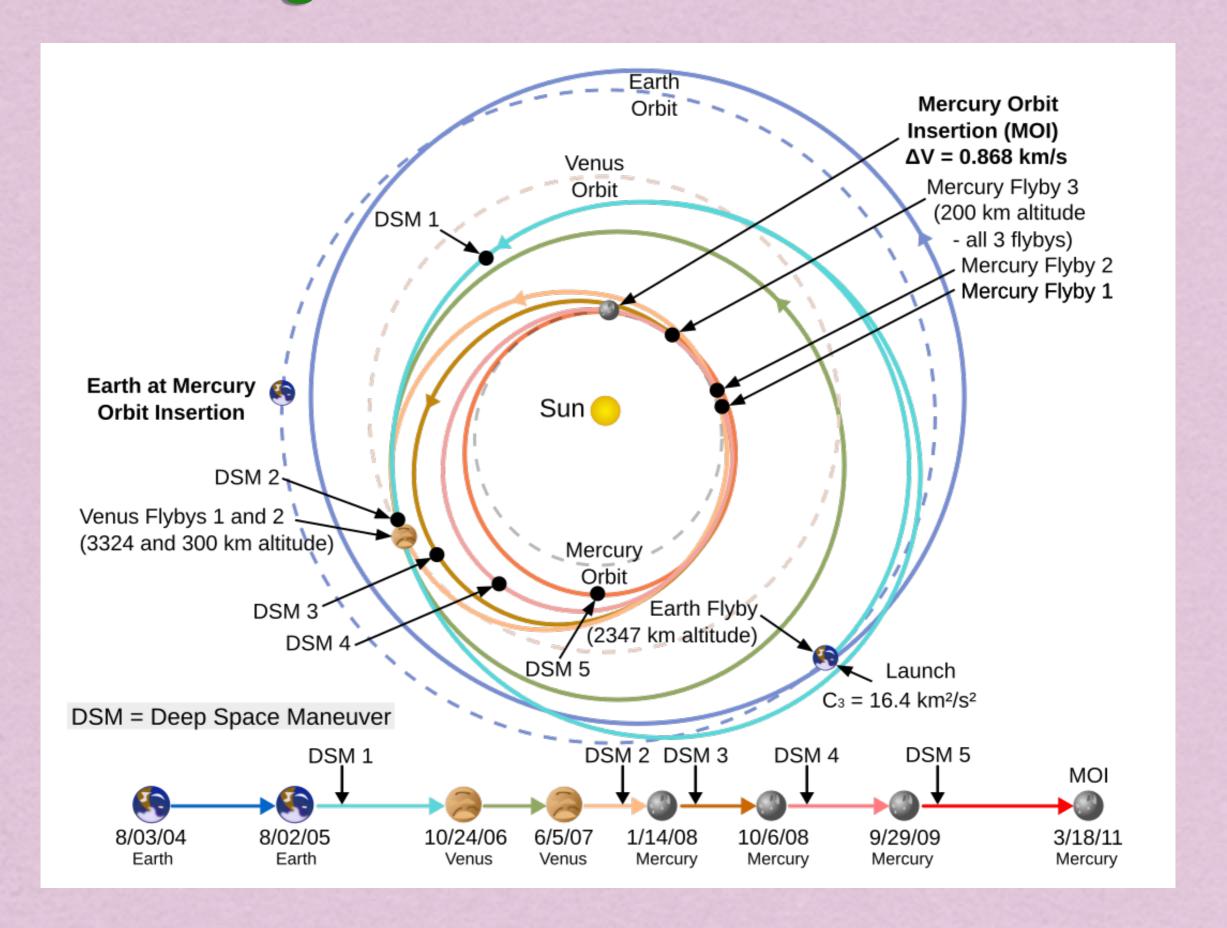
News and Reminders

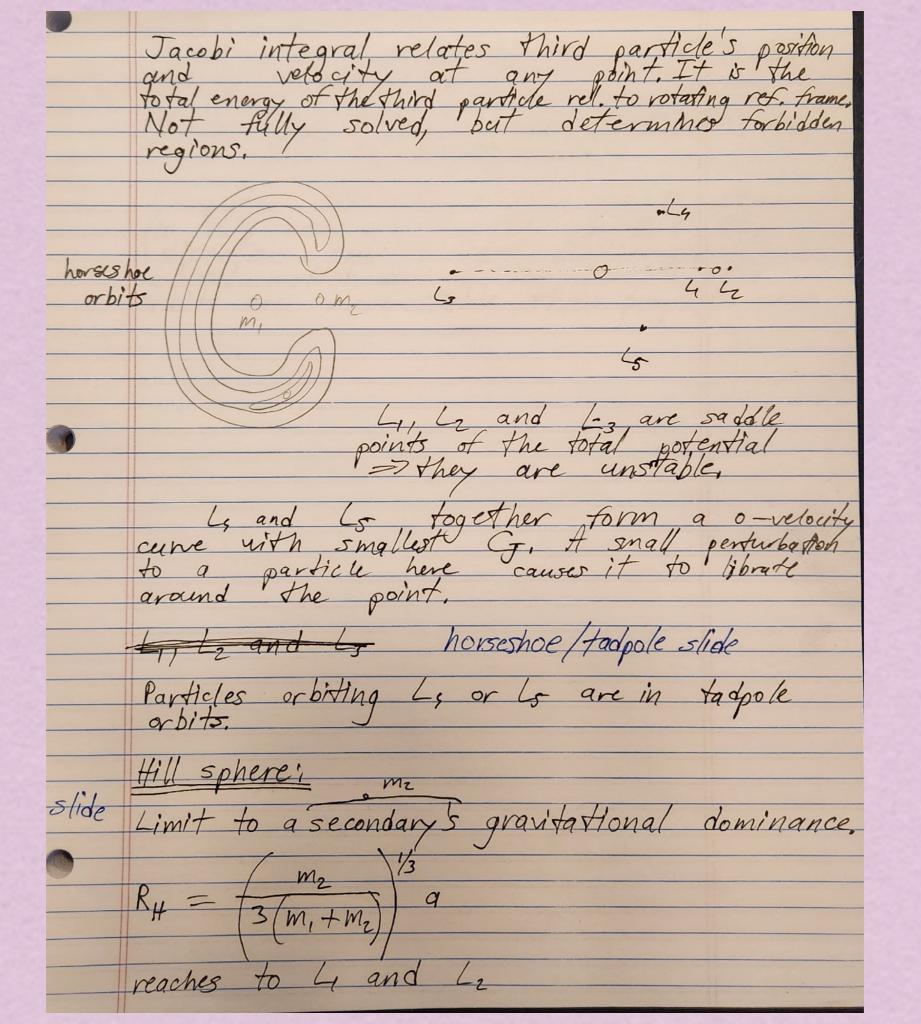
Homework 1 is due now.

Next reading quiz: this Wednesday, 9/11.

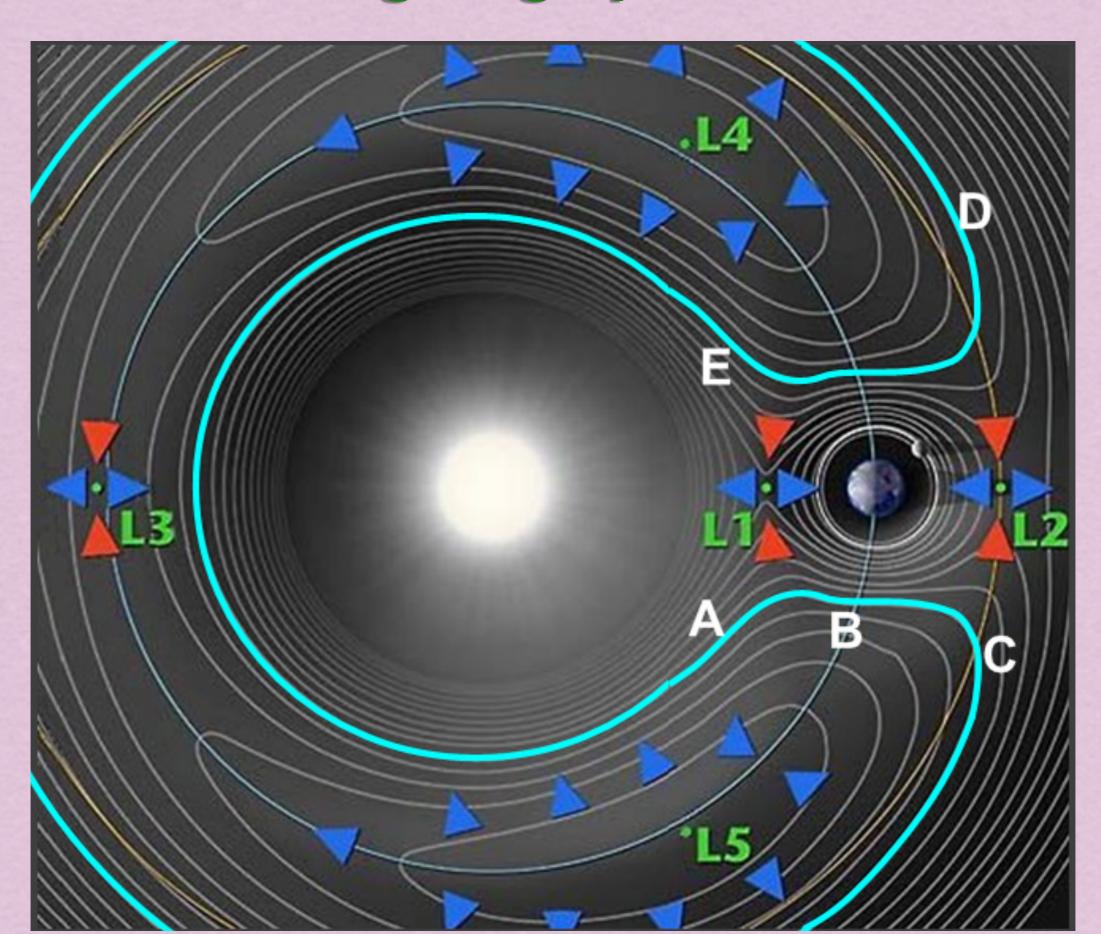
JC 1: Monday, 9/16 (in one week) - Evidence for Hidden Nearby Companions to Hot Jupiters - Sarah Stamer

Messenger Hohmann Transfer Orbits

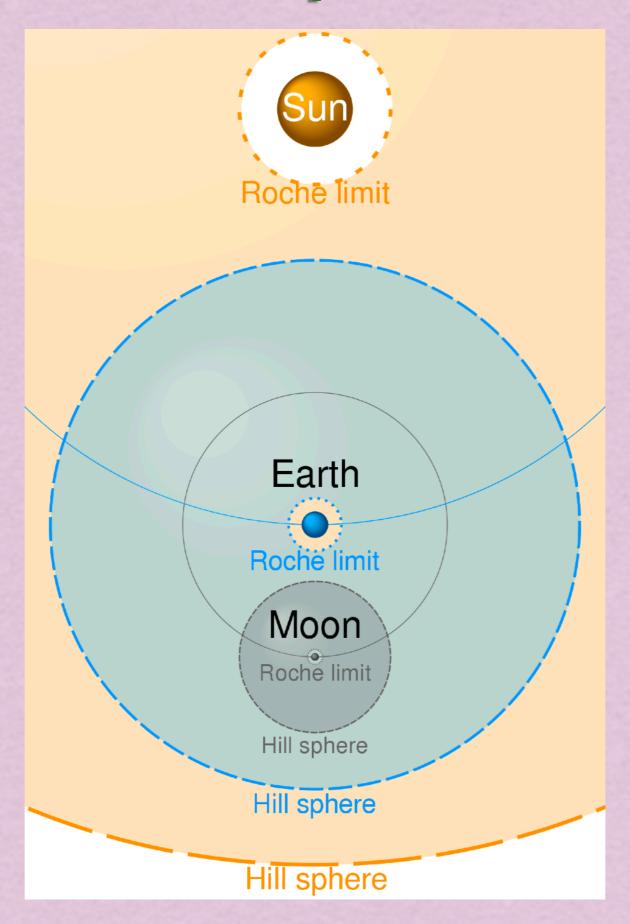




Lagrange points



Hill Sphere

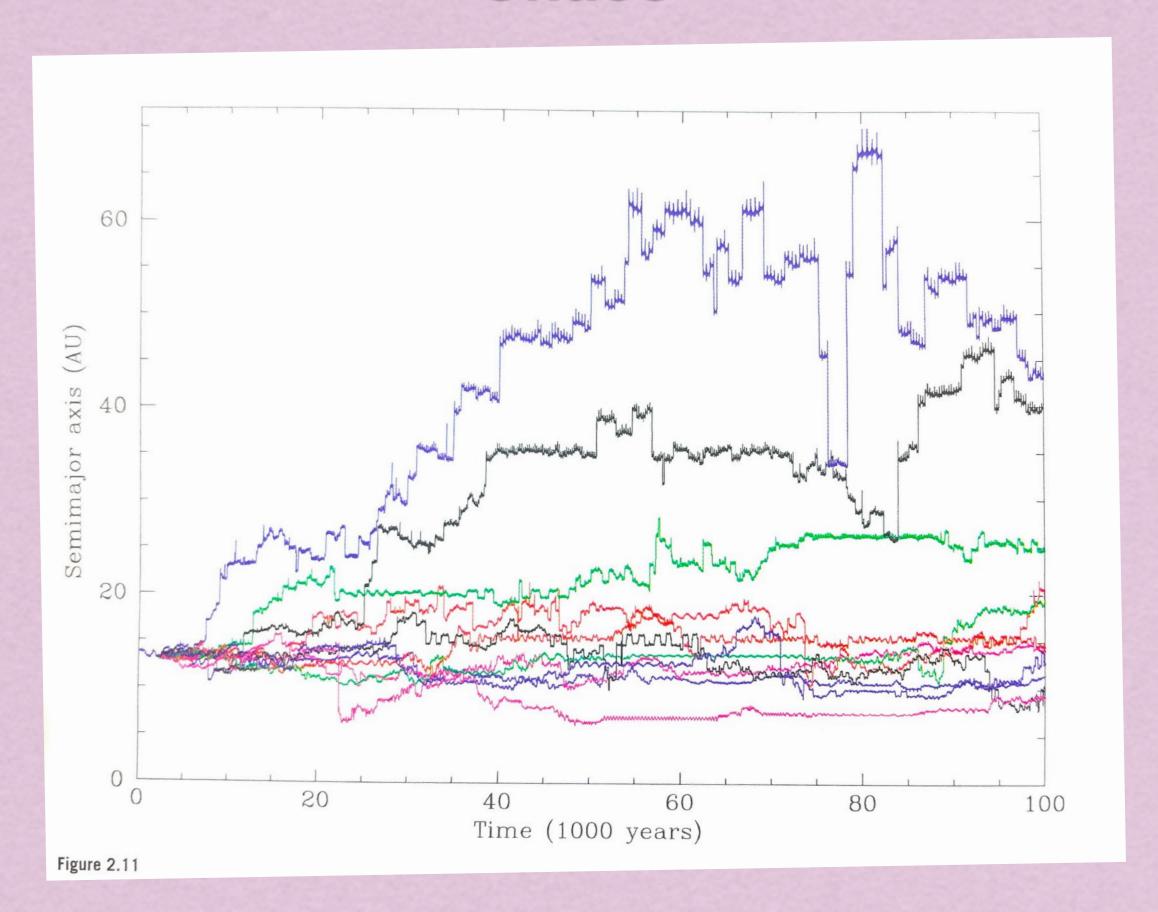


Satellites are well within their planets till sphere, for stability. Objects need to be outside planets HS to be in stable heliocentric orbits. Computer simulations show that retrograde lin this case, orbiting the planet in the direction opposite that of the planets orbital motion around the Sun) satellite orbits are more stable than prograde ones.

Regular VS, Chaotic Motion: We may not be able to analytically solve for motion of orbiting objects most of the time but, at least we can hope to do so numerically. This is true for regular motion; that can be well described nymerically for an arbitrary-long finite time.

So, generally S(t) = S(0) + ctinitial distance bestween distance between two particles two particles at time t But if motion depends too sensitively on the initial conditions, then its behavior becomes effectively unpredictable, and chaos ensues. For chaotic motion: S(t)~ S(o)e" de => Lyapunov exponent/ Je => Lyapunov Hme Class Q: how many Lyapunov times for a 1000 perturbation to result in a 100% discrepancy? Round to nearest integer Interestingly, in the Solar System chaotic behavior develops over timescales orders of magnitude > xet often regions of chaos are associated with resonances ratios of frequencies approximated by rational numbers.

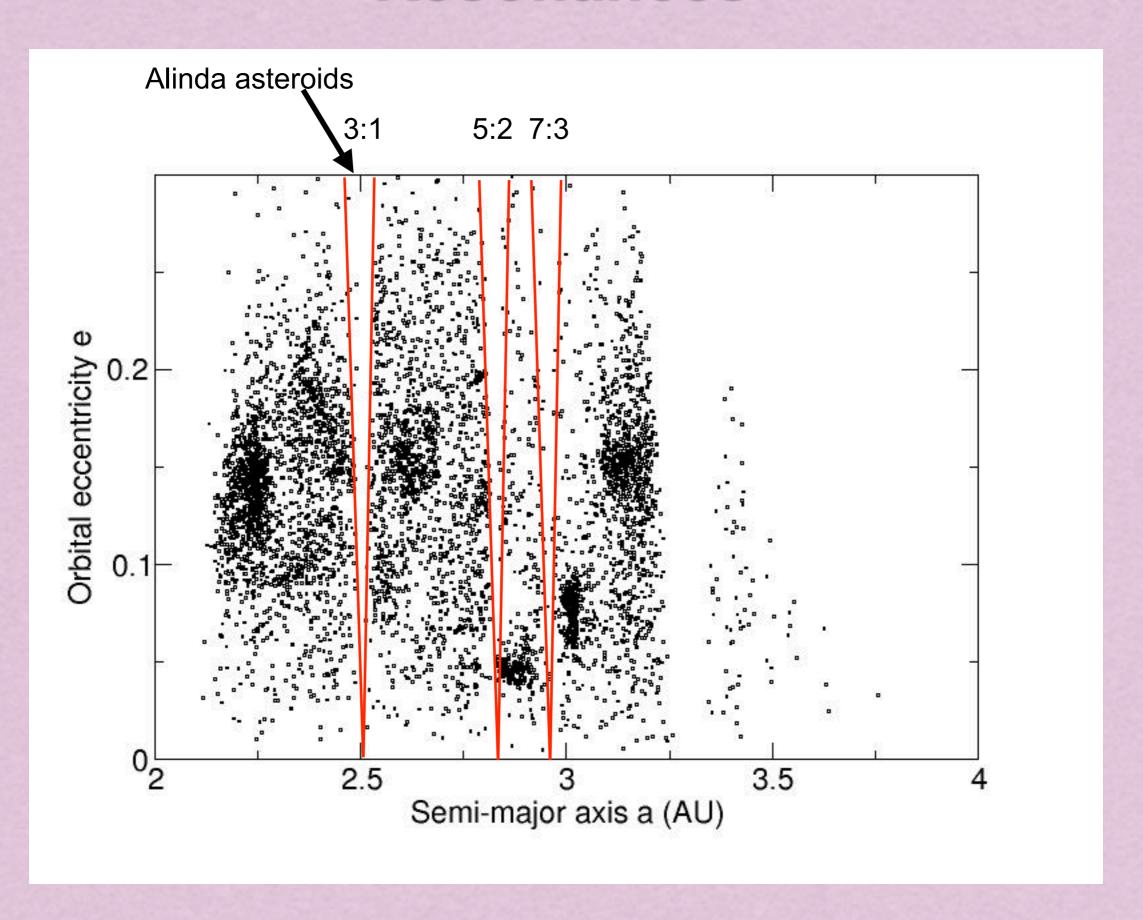
Chaos



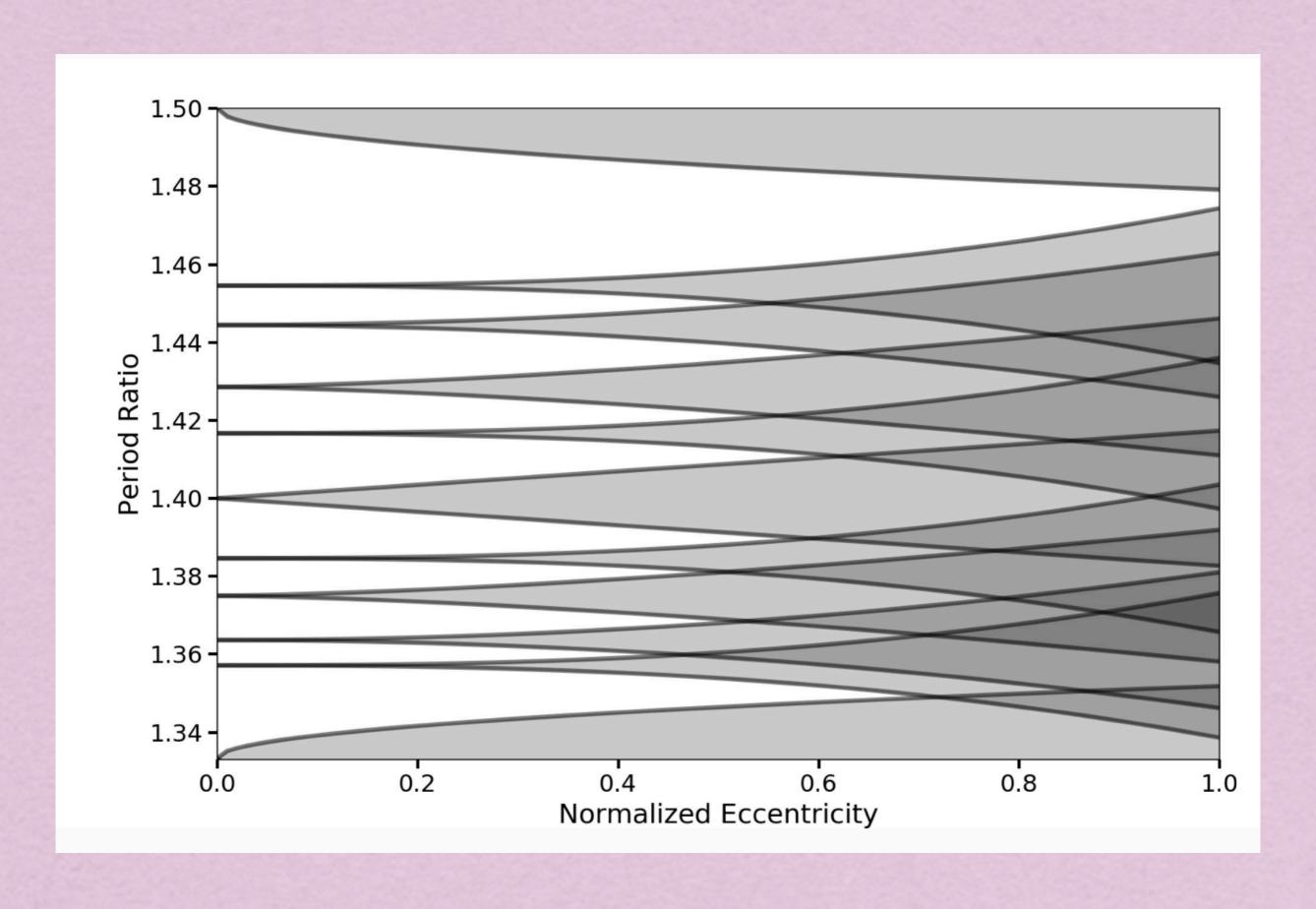
	Wear motion resonance: Att the
	Most long-term perturbations to the motion of a Solar System body are periodic,
	of a Solar System body are periodic,
_	
	sometimes, perturbations even small ones
	Sometimes, perturbations (even small ones) add up coherently => large-amplitude, long-period response
	Cong-period response
	Consider the eggs from of motion of 1D forced
	Consider the equation of motion of 1D forced harmonic oscillator:
	$\frac{md^2x + mw^2x = F_f \cos(w_f t)}{dt^2}$
	Wai natural Frequency
	wf: forging / u/
	wo: natural frequency wf: forcing u integrate to get equation for x:
	Y= Fa cos(u) + + C cos(u) + + C sho(u) +
	$X = \frac{F_f}{m(\omega_0^2 - \omega_f^2)} \frac{\cos(\omega_f t) + G_c \cos(\omega_0 t) + G_s \sin(\omega_0 t)}{\cos(\omega_0 t)}$
	usually small, so x has low amplitude
	Corresponding to initial conditions
	corres ponding to initial conditions.
	What hamens when we & Wo?
	What happens when $w_c \approx w_o$? A large -amplitude response occurs even if Ex is small.
	Fr is small.
	When $w_0 = w_F$, solution is:
	x= Fo + sin (u) + + Coos(u) + + Cos(u) + + Cos(u) +
	2 102(1)
1	x= Ff tsin (wot) + Gcos(wot) + Gz = (wot) 2 mwo 1 steady (secular) growth/

Mean motion resonance resonance applied to orbital period e.g. or N or --higher N: N+1 resonances are stronger because closer to secondary (relative to primary Class: why? but also less stable Resonances in Asteroid belt Class: at what resonance with Jupiter do Stability of the Solar Systemi N-body problem with Lyapunov time: 2-230 Myr => planets or bits are of over long timescales slide e of timescales, Could be: resonance - secular resonance -spin-orbit resonance 5/ide slide on future Mercury and other planets or bits

Resonances

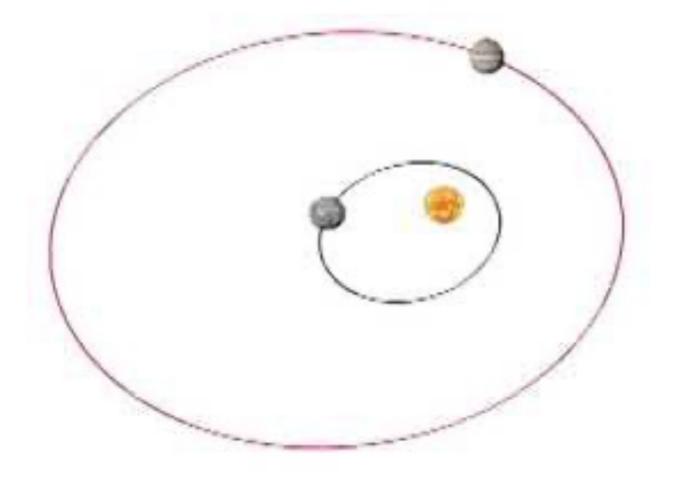


Mean Motion Resonance vs. Eccentricity

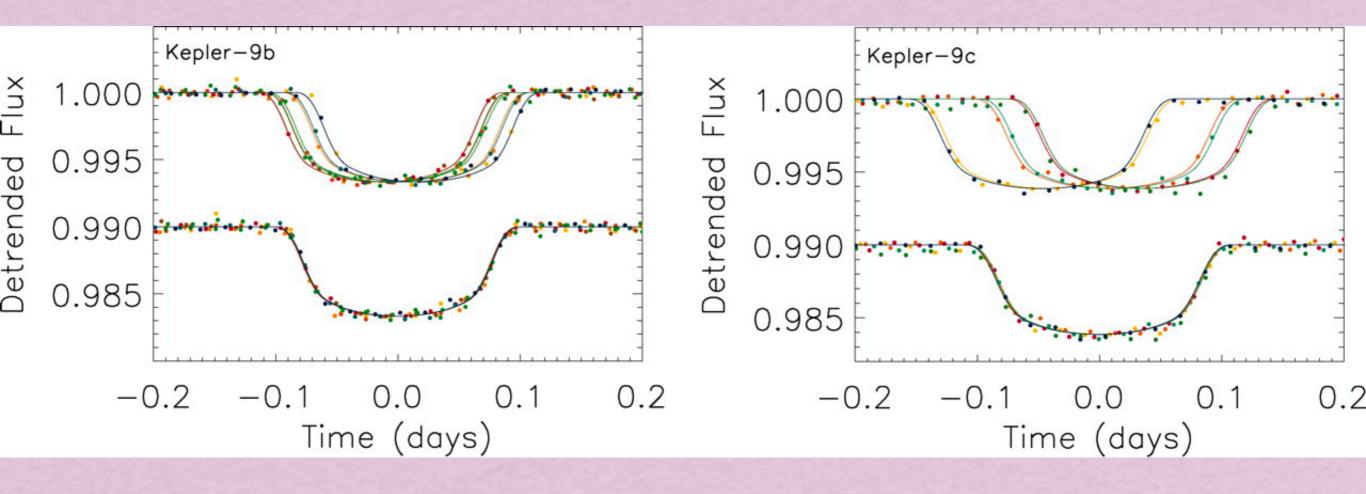


Secular Resonance

Two planets out of resonance



Transit Timing Variations



Holman et al. (2010)

• For Kepler systems, it is often the only way to measure planet masses